

Warped AdS_3 Black Holes: are they classically stable?

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BritGrav 2013, Sheffield



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FCT

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1 Warped AdS_3 black holes

- General properties
- Topologically Massive Gravity
- Warped AdS_3 black holes
- Causal structure of $WAdS_3$ black holes

2 Scalar field perturbation and quasinormal modes

- Scalar field in the background of a $WAdS_3$ black hole
- Quasinormal and bound state modes
- Case with a mirror

3 Conclusions

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- **No** immediate clear notion of time!

WHAT?

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... are they at least classically stable?

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- Technically **simpler!**
- It retains much of the **conceptual complexity** of 4D gravity.
- Einstein gravity equivalent to a **Chern-Simons gauge theory**.

Achucarro, Townsend (1986), Witten (1988,1989)

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with:

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- **Massive** propagating degree of freedom!...
- **Third-order** derivative theory...

Metric of warped AdS_3 black holes

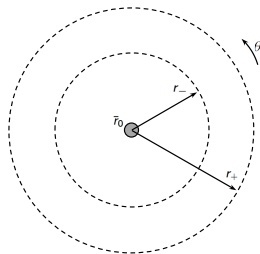
Spacelike stretched black hole:

$$ds^2 = dt^2 + \frac{\ell^2 dr^2}{4R^2(r)N^2(r)} + 2R^2(r)N^\theta(r)dtd\theta + R^2(r)d\theta^2$$

$$R^2(r) = \frac{3(\nu^2 - 1)}{4}r(r - r_0)$$

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Anninos, Li, Padi, Song, Strominger (2008)

$\nu > 1$ is the **warp factor** of the spacetime.

$\nu \rightarrow 1$ recovers the **BTZ black hole**.

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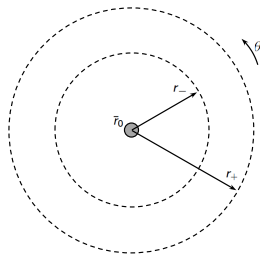
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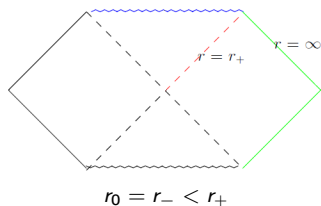
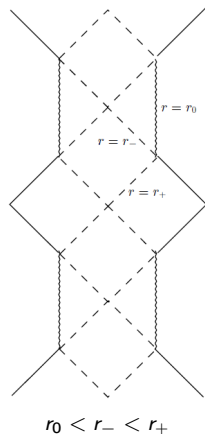
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\implies Observers at radius r follow orbits of $d/dt + \Omega(r) d/d\theta$!

$\Omega(r)$ is restricted to a range of values and:

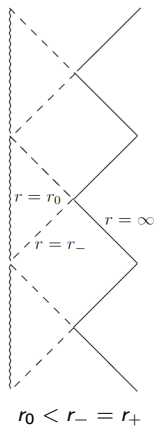
$$\begin{aligned}\Omega(r) &\rightarrow \Omega_{\mathcal{H}}, & r &\rightarrow r_+ \\ \Omega(r) &\rightarrow 0, & r &\rightarrow +\infty\end{aligned}$$

Penrose diagram of a spacelike stretched black hole



- **Not** asymptotically AdS_3 , but similar to asymptotically flat black holes!

- Arena to obtain exact results for a rotating black hole whose causal structure is similar to the Kerr spacetime!



Jugeau, Moutsopoulos, Ritter (2010)

So... is it stable or not?

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Today I will consider **classical** stability.

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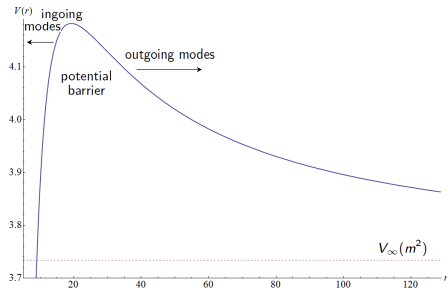
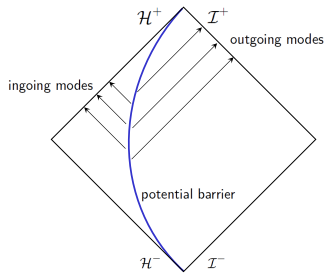
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- **Exact mode solutions:**

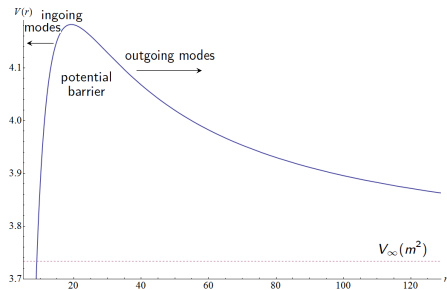
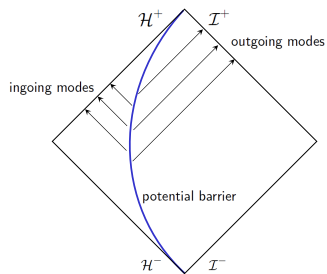
$$\Phi_{\omega k}(t, r, \theta) \sim e^{-i\omega t + ik\theta} z^\alpha (1-z)^\beta F(a, b, c; z)$$

- α, β, a, b, c constants (given in terms of ω and k)
- $z = (r - r_+)/ (r - r_-)$
- $F(a, b, c; z)$ hypergeometric function

What are quasinormal modes?



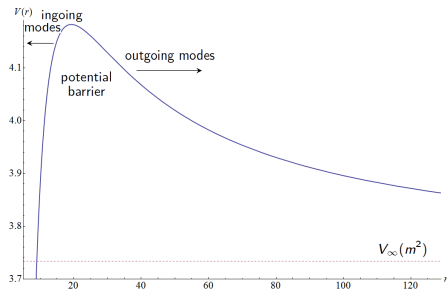
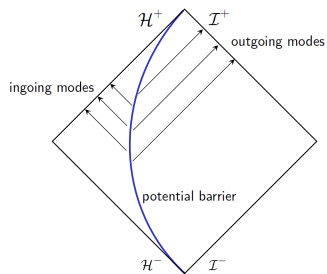
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- **Ingoing** modes at the event horizon;
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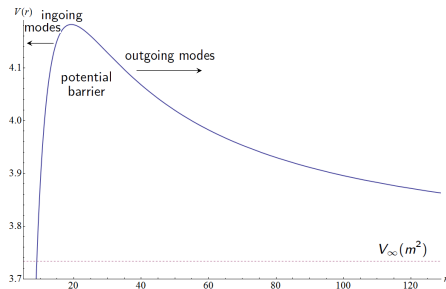
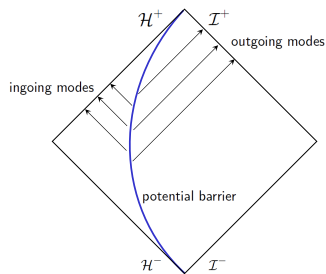
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\implies Discrete set of complex eigenfrequencies $\{\omega_n\}$

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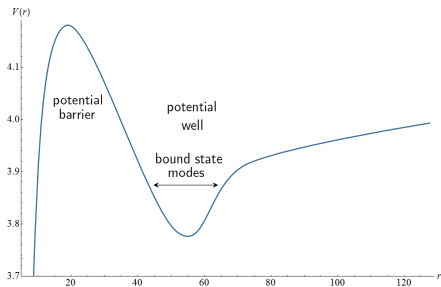
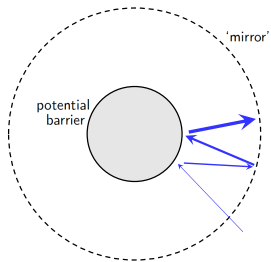
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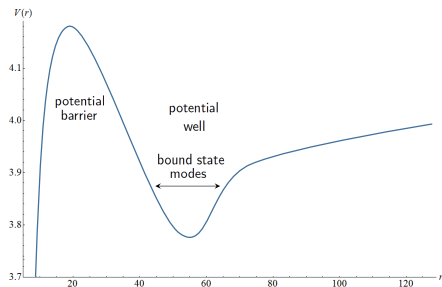
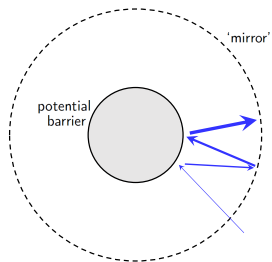
$$\Phi_n \sim e^{-i\omega_n t} = e^{-i\text{Re}(\omega_n)t + \text{Im}(\omega_n)t}$$

If $\text{Im}(\omega_n) > 0$ for some n : mode is **unstable!**

Superradiant instabilities?

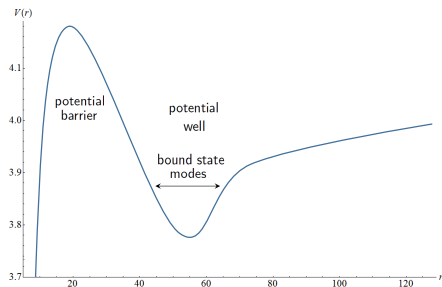
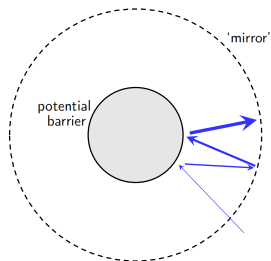


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Bound state modes: superradiant modes that are localised in the potential well (ingoing at event horizon, exponentially decreasing at infinity).

$\implies \text{Im}(\omega_n) > 0 \implies$ **superradiant instability!**

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... NO!

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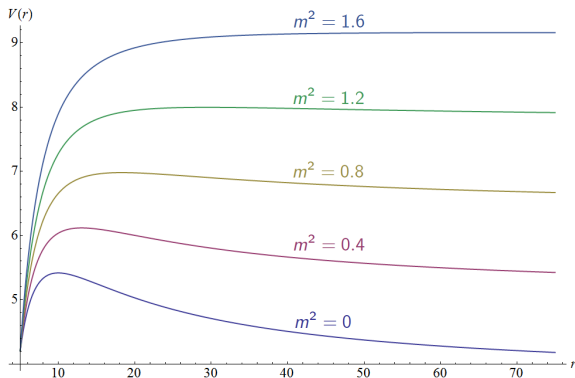
... NO!

- All quasinormal modes have $\text{Im}(\omega_n) < 0$.
- **No** bound state modes!
- **No** superradiant modes in warped AdS_3 black holes (proof?).

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We can obtain exact expressions for the eigenfrequencies!

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+ quasinormal modes - 'exponentially decreasing modes'

■ 'Right eigenfrequencies'

$$(\omega_{\pm})_n^{(R)} = \frac{\nu^2 + 3}{d^2 \delta^2 - 3(\nu^2 - 1)} \left\{ -d\delta \left(\frac{4kd}{\nu^2 + 3} + i \left(n + \frac{1}{2} \right) \right) \pm i(e - i \operatorname{sgn}(k)f) \right\}$$

$$d = \frac{1}{r_+ - r_-}, \quad \delta = 2\nu(r_+ + r_-) - 2\sqrt{(\nu^2 + 3)r_+ r_-},$$

$$e = \sqrt{\frac{\sqrt{E^2 + F^2} + E}{2}}, \quad f = \sqrt{\frac{\sqrt{E^2 + F^2} - E}{2}},$$

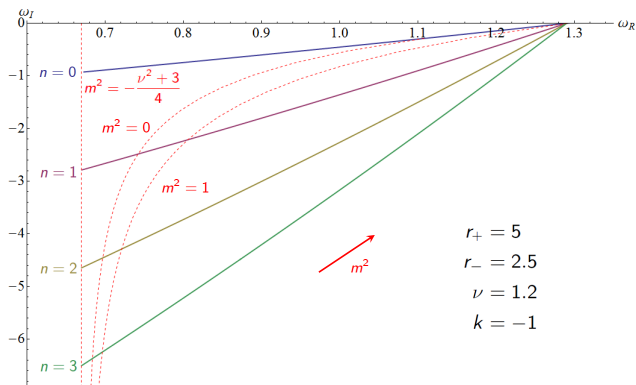
$$E = \left(\frac{1}{4} + \frac{m^2}{\nu^2 + 3} \right) d^2 \delta^2 - 3(\nu^2 - 1) \left[\left(\frac{1}{4} + \frac{m^2}{\nu^2 + 3} \right) + \left(\frac{4kd}{\nu^2 + 3} \right)^2 - \left(n + \frac{1}{2} \right)^2 \right]$$

$$F = -3(\nu^2 - 1) \left(n + \frac{1}{2} \right) \frac{8kd}{\nu^2 + 3}.$$

■ 'Left eigenfrequencies'

$$(\omega_{\pm})_n^{(L)} = -i \left[(2n + 1)\nu \mp \sqrt{3(\nu^2 - 1) \left(n + \frac{1}{2} \right)^2 + (\nu^2 + 3) \left(\frac{1}{4} + \frac{m^2}{\nu^2 + 3} \right)} \right]$$

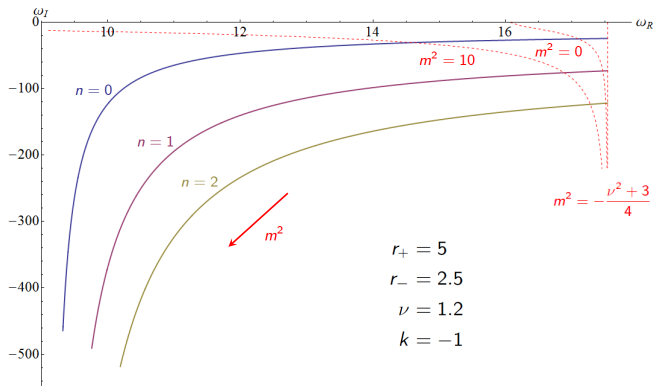
Quasinormal eigenfrequencies in the complex plane w.r.t. m^2



For each n quasinormal modes exist when $m^2 < m_{\max}^2(n)$

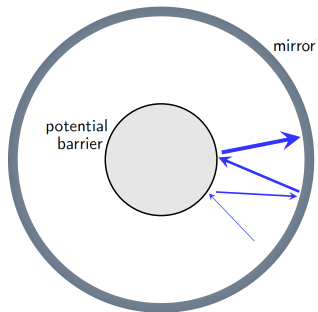
'Exponentially decreasing modes'

'Exponentially decreasing modes' eigenfrequencies w.r.t. m^2

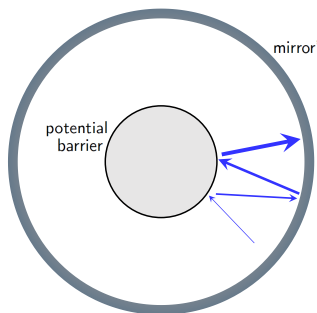


Note: in the context of (W)AdS/CFT, these modes obeying Dirichlet boundary conditions at infinity are *chosen* to be the quasinormal modes.

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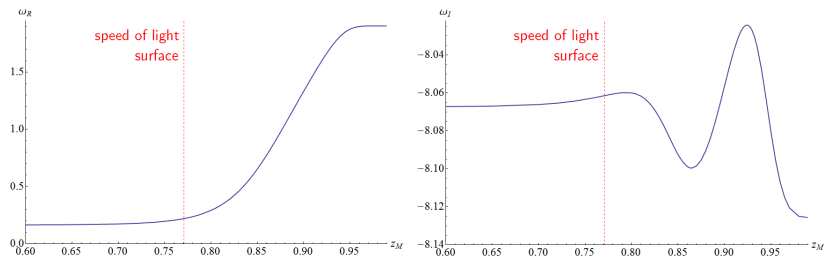


Boundary conditions:

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- **Vanishing** modes at the mirror (Dirichlet boundary condition).

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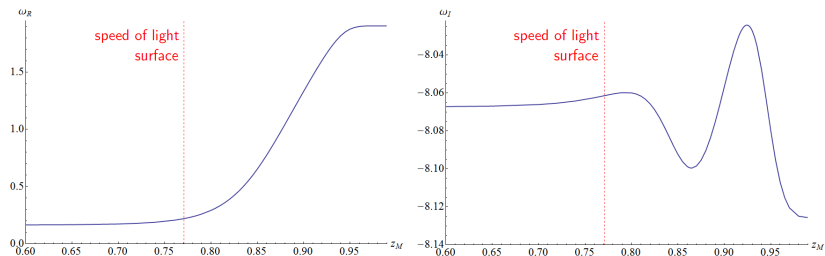
Eigenfrequency vs position of the mirror



$$(r_+ = 7, \quad r_- = 0.7, \quad \nu = 1.2, \quad k = -1, \quad m^2 = 0)$$

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Eigenmodes are still **stable**!

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- Warped AdS_3 black holes has very similar properties to the $(3+1)$ -dimensional Kerr spacetime and, contrary to the latter, many analytical computations can be performed!
- Warped AdS_3 black holes are classically stable to scalar field perturbations, even if a mirror is added to the spacetime.

What's next?

- Is the warped AdS_3 black hole classically stable to other types of perturbations (namely gravitational perturbations)?

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- Is the warped AdS_3 black hole classically stable to other types of perturbations (namely gravitational perturbations)?
- How about quantum effects? Is there a Hartle-Hawking-like vacuum state? What is the renormalised stress-energy tensor for a field in this state?

THANK YOU FOR YOUR ATTENTION!